

$$E_x = A\sqrt{z}J_p(kz) + B\sqrt{z}J_{-p}(kz)$$

where

$$p = [(z_c/\lambda)^2 + 1/4]^{1/2} \geq 1/2.$$

Reference to the asymptotic expression⁶

$$\sqrt{z}J_{-p}(kz) \approx z^{1/2-p},$$

which holds for $z \rightarrow 0$, shows that this part of the solution must be rejected to obtain a finite solution at the origin, corresponding to the case in which the plasma is backed by a metal wall as in Fig. 2(b).⁷ The solution for the plasma with the inverse quadratic density distribution is then

$$E_x = \sqrt{z}J_p(kz)$$

where $p \approx z_c/\lambda$ (assuming $z_c > \lambda$).

Now, comparing (6) and (10) (with $\theta=0$) we see that the waveguide and plasma shown in Fig. 2 exhibit analogous behavior with respect to a TE wave propagating toward the origin. In particular, the wave undergoes "reflection" in the neighborhood of the critical density z_c or the critical cross section r_c . Similarly, an evanescent wave arises at the critical point and damps out toward the origin. Thus, analog experiments carried out with appropriate tapered-waveguide configurations appear to be useful for simulating inhomogeneous plasma media in the vicinity of critical densities. Such experiments would require slowly-varying taper structures in order to minimize effects due to generation of spurious modes at junctions with other microwave elements.

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⁷ L. S. Taylor, "Reflection of a TE wave from an inverse parabolic ionization density," IRE TRANS. ON ANTENNAS AND PROPAGATION (Correspondence), vol. AP-9, pp. 582-583; November, 1961.

The Diffraction Loss Curve for Nonconfocal Spherical Mirrors

The diffraction loss curve vs Fresnel number for confocal spherical mirrors was shown by Fox and Li¹ and by Goubau and Schwering.² Boyd and Gordon³ suggested a possibility of applying the above theory to a nonconfocal mirror system, by assuming the diffraction loss to be equal to that of an equivalent confocal system having the same

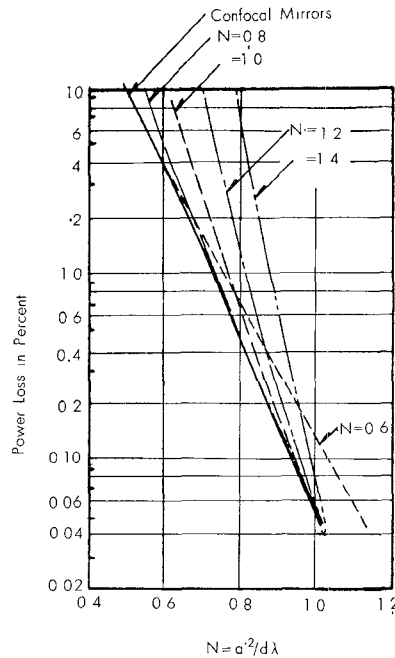


Fig. 1—Diffraction loss for a nonconfocal system of spherical mirrors.

spot size. They proposed to express the diffraction loss in term of the parameter

$$\frac{a'^2}{d\lambda} \left[2 \frac{d}{b'} - \left(\frac{d}{b'} \right)^2 \right]^{1/2}$$

in place of the conventional Fresnel number N .

The approximation by Boyd and Gordon has been found nearly valid by Fox and Li¹ for the range of $0.2b' < d < 1.8b'$ in the calculation of infinite strip curved mirror interferometers for $N=0.5$.

Using the Boyd and Gordon approximation, the diffraction loss for a nonconfocal system may be expressed as a function of the above parameter, which can be modified to the form

$$N \left[2 \left(\frac{b'}{d} \right) - 1 \right]^{1/2}$$

where

b' = radius of curvature of mirrors

d = spacing between mirrors

N = Fresnel number

$= a'^2/b'\lambda$

a' = radius of mirrors.

The above new parameter corresponds to the Fresnel number N for a confocal system, and the diffraction loss for a nonconfocal system can be easily obtained using the loss curve for a confocal system by replacing N to the form modified by the factor

$$\left[2 \left(\frac{b'}{d} \right) - 1 \right]^{1/2}.$$

On the other hand, it is sometimes required to calculate the variation of the diffraction loss for a nonconfocal system with

⁴ A. G. Fox and T. Li, "Modes in a maser interferometer with curved and tilted mirrors," PROC. IEEE, vol. 51, pp. 80-89; January, 1963.

the spacing between mirrors for constant mirror curvature and wavelength. In such a case, it seems more convenient to use the new parameter N' as defined by the formula

$$N' = \frac{a'^2}{d\lambda}.$$

With this parameter, the diffraction loss for a nonconfocal system can be illustrated for various values of N as shown in Fig. 1. The solid line in the figure denotes the loss curve for a confocal system for comparison. These curves seem to have sufficient accuracy in the range of $0.2b' < d < 1.8b'$.

When the diffraction loss for a nonconfocal system was measured for a variable mirror spacing, the measured value should be compared with the curve shown in Fig. 1, not with the loss curve for a confocal system. It is considered that the results of the measurement made by Beyer and Scheibe⁵ may be compared more adequately with the curve in Fig. 1 for a given value of N .

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⁵ J. B. Beyer and E. H. Scheibe, "Loss measurements of the beam waveguide," IEEE TRANS. ON MICROWAVE THEORY AND TECHNIQUES, vol. MTT-11, pp. 18-22; January, 1963.

Semiconductor Switching and Limiting Using 3-db Short-Slot (Hybrid) Couplers

The silver-bonded germanium varactor diode has been successfully used as a switch and a limiter of microwave power when operated in a series mode between 9.0 and 9.6 Gc.¹ This report gives details of shunt mode switching and limiting using these same type diodes in conjunction with 3-db short-slot (hybrid) couplers. The technique of using 3-db short-slot (hybrid) couplers, but with other type diodes (e.g., 1N263, MA-450, PIN's), has been reported by other investigators.²⁻⁴

Fig. 1 is a diagrammatic illustration of the short-slot (hybrid) coupler. If arms B and C are terminated in perfectly matched

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¹ V. J. Higgins, "X-band semiconductor switching and limiting using waveguide series tees," Microwave J., vol. 6, pp. 77-83; November, 1963.

² R. Lucy, "Microwave High Speed Switch," Proc. Natl. Electronics Components Conf., Philadelphia, Pa., pp. 12-15; May, 1959.

³ R. V. Garver and D. V. Tseng, "X-band diode limiting," IRE TRANS. ON MICROWAVE THEORY AND TECHNIQUES (Correspondence), vol. MTT-9, p. 202; March, 1961.

⁴ W. F. Krupke, T. S. Hartwick and M. T. Weiss, "Solid-state X-band power limiter," IRE TRANS. ON MICROWAVE THEORY AND TECHNIQUES, vol. MTT-9, pp. 472-480; November, 1961.

Manuscript received December 17, 1963.

¹ A. G. Fox and T. Li, "Resonant modes in a maser interferometer," Bell Sys. Tech. J., vol. 40, pp. 453-488; March, 1961.

² G. Goubau and F. Schwering, "On the guided propagation of electromagnetic wave beams," IRE TRANS. ON ANTENNAS AND PROPAGATION, pp. 248-256; May, 1961.

³ G. D. Boyd and J. P. Gordon, "Confocal multimode resonator for millimeter through optical wavelength masers," Bell Sys. Tech. J., vol. 40, pp. 489-508; March, 1961.