

$$E_x = A\sqrt{z}J_p(kz) + B\sqrt{z}J_{-p}(kz)$$

where

$$p = [(z_c/\lambda)^2 + 1/4]^{1/2} \geq 1/2.$$

Reference to the asymptotic expression<sup>6</sup>

$$\sqrt{z}J_{-p}(kz) \approx z^{1/2-p},$$

which holds for  $z \rightarrow 0$ , shows that this part of the solution must be rejected to obtain a finite solution at the origin, corresponding to the case in which the plasma is backed by a metal wall as in Fig. 2(b).<sup>7</sup> The solution for the plasma with the inverse quadratic density distribution is then

$$E_x = \sqrt{z}J_p(kz)$$

where  $p \approx z_c/\lambda$  (assuming  $z_c > \lambda$ ).

Now, comparing (6) and (10) (with  $\theta = 0$ ) we see that the waveguide and plasma shown in Fig. 2 exhibit analogous behavior with respect to a TE wave propagating toward the origin. In particular, the wave undergoes "reflection" in the neighborhood of the critical density  $z_c$  or the critical cross section  $r_c$ . Similarly, an evanescent wave arises at the critical point and damps out toward the origin. Thus, analog experiments carried out with appropriate tapered-waveguide configurations appear to be useful for simulating inhomogeneous plasma media in the vicinity of critical densities. Such experiments would require slowly-varying taper structures in order to minimize effects due to generation of spurious modes at junctions with other microwave elements.

HERBERT LASHINSKY  
Plasma Physics Lab.  
Princeton University  
Princeton, N. J.

<sup>7</sup> L. S. Taylor, "Reflection of a TE wave from an inverse parabolic ionization density," *IRE TRANS. ON ANTENNAS AND PROPAGATION (Correspondence)*, vol. AP-9, pp. 582-583; November, 1961.

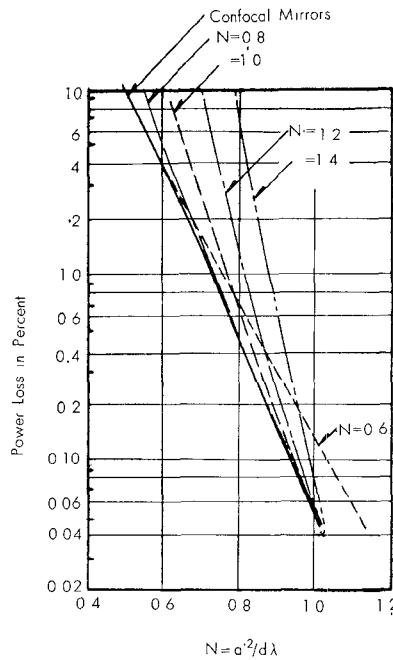


Fig. 1—Diffraction loss for a nonconfocal system of spherical mirrors.

spot size. They proposed to express the diffraction loss in term of the parameter

$$\frac{a'^2}{d\lambda} \left[ 2 \frac{d}{b'} - \left( \frac{d}{b'} \right)^2 \right]^{1/2}$$

in place of the conventional Fresnel number  $N$ .

The approximation by Boyd and Gordon has been found nearly valid by Fox and Li<sup>1</sup> for the range of  $0.2b' < d < 1.8b'$  in the calculation of infinite strip curved mirror interferometers for  $N = 0.5$ .

Using the Boyd and Gordon approximation, the diffraction loss for a nonconfocal system may be expressed as a function of the above parameter, which can be modified to the form

$$N \left[ 2 \left( \frac{b'}{d} \right) - 1 \right]^{1/2}$$

where

$b'$  = radius of curvature of mirrors

$d$  = spacing between mirrors

$N$  = Fresnel number

$= a'^2/b'^2$

$a'$  = radius of mirrors.

The above new parameter corresponds to the Fresnel number  $N$  for a confocal system, and the diffraction loss for a nonconfocal system can be easily obtained using the loss curve for a confocal system by replacing  $N$  to the form modified by the factor

$$\left[ 2 \left( \frac{b'}{d} \right) - 1 \right]^{1/2}.$$

On the other hand, it is sometimes required to calculate the variation of the diffraction loss for a nonconfocal system with

<sup>1</sup> A. G. Fox and T. Li, "Resonant modes in a maser interferometer," *Bell Syst. Tech. J.*, vol. 40, pp. 453-488; March, 1961.

<sup>2</sup> G. Goubau and F. Schwering, "On the guided propagation of electromagnetic wave beams," *IRE TRANS. ON ANTENNAS AND PROPAGATION*, pp. 248-256; May, 1961.

<sup>3</sup> G. D. Boyd and J. P. Gordon, "Confocal multimode resonator for millimeter through optical wavelength masers," *Bell Syst. Tech. J.*, vol. 40, pp. 489-508; March, 1961.

the spacing between mirrors for constant mirror curvature and wavelength. In such a case, it seems more convenient to use the new parameter  $N'$  as defined by the formula

$$N' = \frac{a'^2}{d\lambda}.$$

With this parameter, the diffraction loss for a nonconfocal system can be illustrated for various values of  $N$  as shown in Fig. 1. The solid line in the figure denotes the loss curve for a confocal system for comparison. These curves seem to have sufficient accuracy in the range of  $0.2b' < d < 1.8b'$ .

When the diffraction loss for a nonconfocal system was measured for a variable mirror spacing, the measured value should be compared with the curve shown in Fig. 1, not with the loss curve for a confocal system. It is considered that the results of the measurement made by Beyer and Scheibe<sup>8</sup> may be compared more adequately with the curve in Fig. 1 for a given value of  $N$ .

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TAKEO OMORI  
Electrical Communication Lab.  
Nippon Telephone and Telegraph  
Public Corp.  
Tokyo, Japan

<sup>8</sup> J. B. Beyer and E. H. Scheibe, "Loss measurements of the beam waveguide," *IEEE TRANS. ON MICROWAVE THEORY AND TECHNIQUES*, vol. MTT-11, pp. 18-22; January, 1963.

#### Semiconductor Switching and Limiting Using 3-db Short-Slot (Hybrid) Couplers

The silver-bonded germanium varactor diode has been successfully used as a switch and a limiter of microwave power when operated in a series mode between 9.0 and 9.6 Gc.<sup>1</sup> This report gives details of shunt mode switching and limiting using these same type diodes in conjunction with 3-db short-slot (hybrid) couplers. The technique of using 3-db short-slot (hybrid) couplers, but with other type diodes (e.g., 1N263, MA-450, PIN's), has been reported by other investigators.<sup>2-4</sup>

Fig. 1 is a diagrammatic illustration of the short-slot (hybrid) coupler. If arms  $B$  and  $C$  are terminated in perfectly matched

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<sup>1</sup> W. J. Higgins, "X-band semiconductor switching and limiting using waveguide series tees," *Microwave J.*, vol. 6, pp. 77-83; November, 1963.

<sup>2</sup> R. Lucy, "Microwave High Speed Switch," *Proc. Natl. Electronics Components Conf.*, Philadelphia, Pa., pp. 12-15; May, 1959.

<sup>3</sup> R. V. Garver and D. V. Tseng, "X-band diode limiting," *IRE TRANS. ON MICROWAVE THEORY AND TECHNIQUES (Correspondence)*, vol. MTT-9, p. 202; March, 1961.

<sup>4</sup> W. F. Krupke, T. S. Hartwick and M. T. Weiss, "Solid-state X-band power limiter," *IRE TRANS. ON MICROWAVE THEORY AND TECHNIQUES*, vol. MTT-9, pp. 472-480; November, 1961.